

SPECTROSCOPIC MEASUREMENTS OF NON-MAXWELLIAN AND ANISOTROPIC FEATURES OF THE ION DISTRIBUTION FUNCTION IN H-1 HELIAC



C.A. Michael, J. Howard
PRL,RSPHysSE, ANU, Canberra A.C.T. 0200 Australia

COHERENCE MEASUREMENTS WITH MOSS

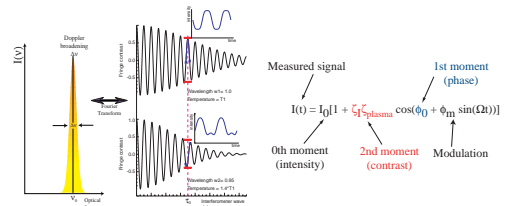
The Modulated optical Solid State spectrometer (MOSS) is a fixed delay optical Fourier transform spectrometer. It can measure the light intensity, spectral width and centre frequency, by monitoring the coherence (fringe contrast and phase) at one fixed delay. The centre frequency (phase) conveys the ion flow speed, and, assuming that the distribution function is Maxwellian, a temperature can be determined from the contrast.

The measured contrast and phase is related to the lineshape by the Fourier transform:

$$\zeta(\tau)e^{i\phi(\tau)} = \Gamma(\tau) = F(I(v))$$

contrast
coherence
lineshape
phase
(aka distribution function)

The following figures illustrate the measurement principle:

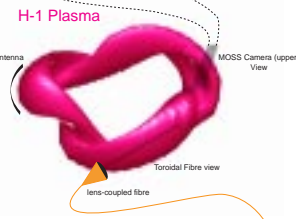


MOSS DIAGNOSTICS ON H-1



MOSS CAMERA:
 * Channels: 32
 * Coverage: Entire plasma cross section
 * Direct View
 * Waveplates can be added between shots

The MOSS spectrometer is a pocket cell, consisting of an electro-optic, birefringent, modulated wave-plate, polarisers and an interference filter to select the line of interest [1].



THE MOSS CAMERA images a poloidal cross-sectional view of the plasma, as shown below. This instrument is sensitive to the temperature and flow mostly **PERPENDICULAR** to the magnetic field.

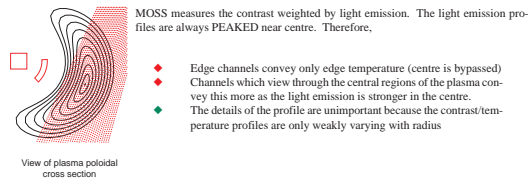
A **FIBRE COUPLED MOSS** system views a single chord running through the entire plasma cross section, in a direction mostly **PARALLEL** to the magnetic field.

With the two diagnostics operated together, it's possible to measure the **ISOTROPY** of the distribution function, by comparing poloidal and toroidal temperatures.



FIBRE-COUPLED SINGLE-CHANNEL DUAL MOSS
 The incoming light is separated into two separate waveplates, each of different delay, to check the distribution function.
 * Waveplates: fixed (not changed between shots)

LINE-INTEGRATION EFFECT



ABSTRACT

In this poster, we show recent results from coherence imaging spectroscopic systems on H-1, examining the ion distribution function in low-field (B<0.15T) Argon discharges (recall poster S80 on Monday by J. Howard for further information). Strong non-Maxwellian features and anisotropic characteristics have been found. Specifically, most of the measured distribution functions can be fitted to a "two temperature" Maxwellian incorporating both a hot and minority (20%) cold population. The cold component of the distribution function has a larger proportion at the edge, which is consistent with the explanation that it is caused by charge exchange with neutrals. The parallel ion temperature is smaller than the perpendicular component at lower background fill pressures and at higher magnetic fields, which may indicate that the heating mechanism drives the perpendicular component preferentially.

EXPERIMENT

Measurements of contrast were made with the MOSS camera at different crystal thicknesses, in order to test whether the distribution function was Maxwellian.

Function: $\text{contrast } (\zeta) \xleftarrow{F \text{ (fourier transform)}} \text{distribution function } (f)$
 Variable: $\text{crystal thickness } (L) \xrightarrow{\text{velocity } (v)}$

Crystals were interchanged in between shots, so that by relying on shot-to-shot reproducibility, a spatial map of the distribution function could be made.

For a Maxwellian distribution, $f(v)$ is a Gaussian. Therefore $\zeta(L)$ is gaussian.

However, in some instances, a gaussian fits the data poorly. Here a "two temperature" Maxwellian was fitted:

$$\zeta = e^{-\frac{L^2}{w^2}} \text{ Single Gaussian}$$

$$\zeta = (1 - \alpha)e^{-\frac{L^2}{w_1^2}} + \alpha e^{-\frac{L^2}{w_2^2}} \text{ Double Gaussian (two temperatures)}$$

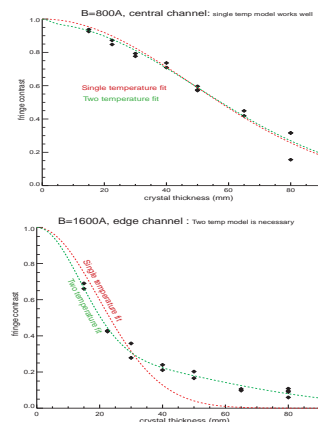
The "characteristic widths" can be converted to temperatures, according to:

$$T_i = \frac{\lambda^2 m_i^2 c^3}{8k_B} w_i^{-2}$$

Therefore,
 * **HOT IONS** ==> SMALL LENGTH SCALE
 * **COLD IONS** ==> LARGE LENGTH SCALE

Most of the data show a "Tail" in the coherence curve $\zeta(L)$, implying the existence of a cold component.

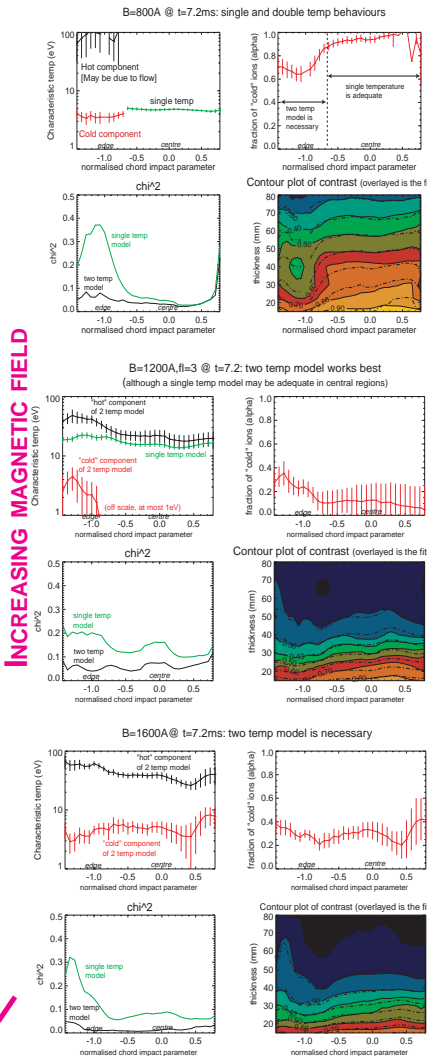
The applicability of each model is shown by the following data:



NON-MAXWELLIAN FEATURES

Shown below are a series of graphs showing profiles of temperature, the chi-squared (residual error) of the fit to each model and the contrast. The applicability a single or double gaussian fit can be judged by:

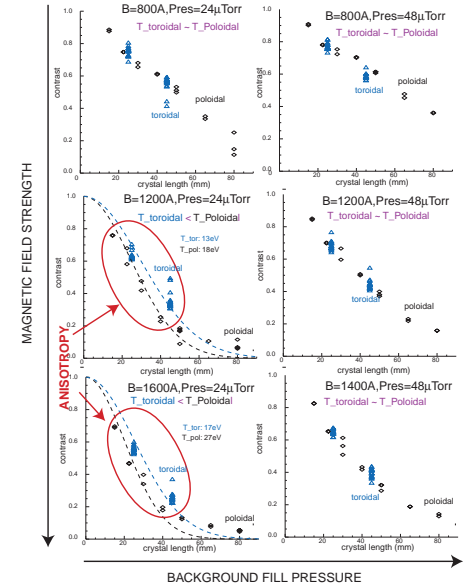
- * Looking at the fraction of cold/hot ions "alpha" (close to 1 or 0 implies that a single temperature description is adequate)
- * Examining the chi-squared residual. When the chi-squared for a two-temperature model is not much better than that of a one-temperature model, a single temperature is all that is necessary to characterise the distribution. Note that the chi-squared does not go to zero because of the scatter in the data (shot-to-shot reproducibility)



ANISOTROPY

The isotropy of the distribution function is measured by comparing the contrast curves (temperatures) from the toroidal viewing MOSS with a spatial average of the poloidal viewing MOSS.

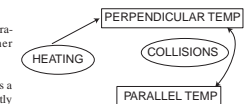
Shown below is the contrast as a function of crystal thickness, at different plasma conditions (filling pressure and magnetic field). The distribution is isotropic where the blue points lie amongst the black ones, however, at high field and low pressure, an anisotropy is detected, and fits to a single temperature model are indicated.



INTERPRETATION

ANISOTROPY

- * Parallel and perpendicular temperatures will equilibrate at higher collisionality (pressure).
- * The existence of anisotropy implies a heating mechanism which dominantly drives the perpendicular component.



TWO-TEMPERATURE DISTRIBUTION

- * Charge-exchange collisions may effect the ions, because of the large background of neutrals in H-1. This will drain some of the energy from ions, creating a "cold" component. This effect should be more pronounced at the edge, where there is an influx of neutrals. Measurements show a larger cold component at the edge.

ACKNOWLEDGEMENTS

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REFERENCES

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 MOSS Camera: Rev Sci. Instrum, vol.72, 2001, pp.1034-7